

Figure 1. Differential scanning calorimetry trace of the triblock polyurethanes.

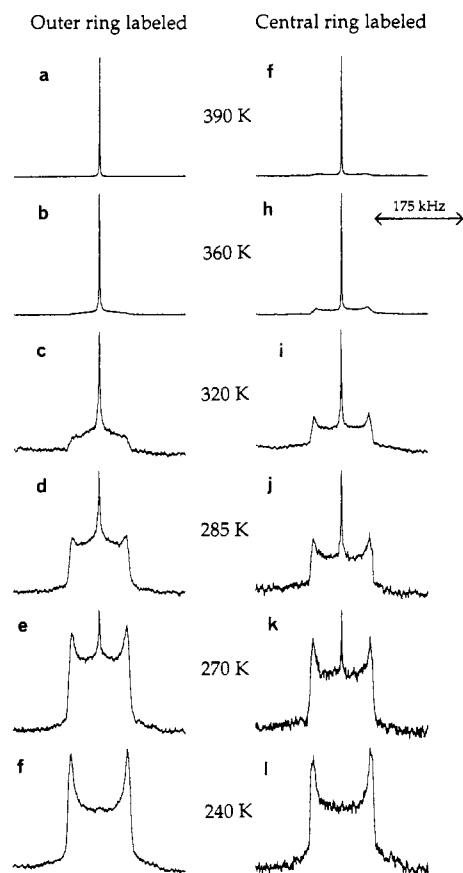


Figure 2. ^2H NMR spectra (46.07 MHz) of the piperazine-based triblock polyurethanes, selectively labeled in the hard segment. Spectra a–f were obtained from the sample labeled in the outer rings, and spectra g–l were obtained from the sample labeled in the central ring.

frequency angular fluctuations, also observed in all hard-segment materials of MDI-BDO-based polyurethanes⁶ and in other crystalline polymers.¹³ Furthermore, the spectra of both the outer- and center-labeled triblock system are dominated by the narrow, motionally averaged peak (line width ca. 1 kHz at 300 K) at temperatures well below $T_m(\text{hard})$. In fact, for $410\text{ K} < T < 440\text{ K}$, the broad spectral features are not detected at all. From a comparison of melting enthalpies for the hard domains in the triblock polyurethanes and the crystalline hard-segment oligomer⁹ an "overall" crystallinity of at least 70% is estimated. Thus, at most a 30% motionally narrowed signal can result from noncrystalline hard segments at temperatures above their T_g , estimated to be around $313 \pm 10\text{ K}$. The temperature dependence of the fraction of

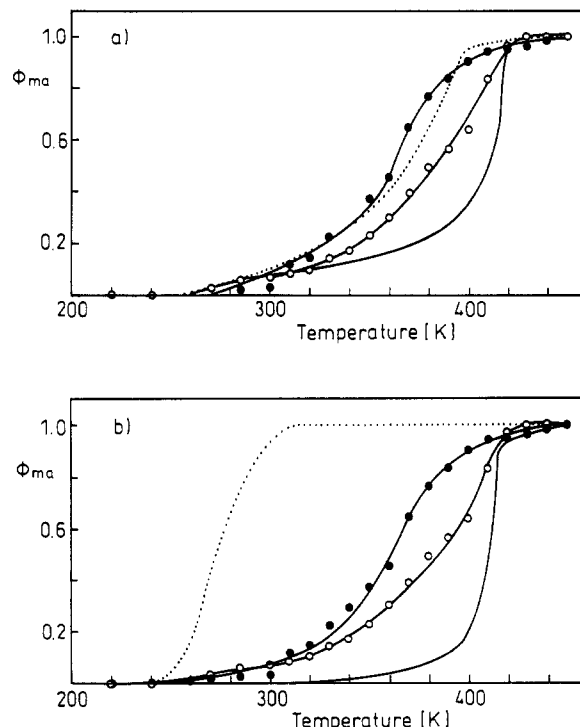


Figure 3. Fraction motionally averaged, Φ_{ma} , of the polyurethane triblocks: (●) outer-labeled triblock polyurethane; (○) center-labeled triblock polyurethane. (a) (—) Center-labeled multiblock polyurethane (reproduced from ref 9); (···) outer-labeled multiblock polyurethane (reproduced from ref 9). (b) (—) Center-labeled hard-segment oligomer (reproduced from ref 9); (···) polyurethane that is chain extended with individual piperazine rings (reproduced from ref 9).

labels undergoing rapid motions for whose NMR signal is motionally averaged, Φ_{ma} , is plotted in Figure 3. For the center-labeled system (see Figure 3a), this limit is reached around 360 K already. Thus the narrow component must in part be attributed to hard segments undergoing quasi-isotropic motion within the crystallites, at least at elevated temperatures.

This quasi-isotropic motion with $\tau_c < 10^{-7}\text{ s}$ may involve discrete jumps or diffusive motions or a combination of both. X-ray scattering investigations into the molecular structure and packing of the piperazine rings indicate that the angles formed by the D–C–D bonds do not deviate to a significant extent from the true tetrahedral angle,¹⁴ and, hence, effective tetrahedral jumps are likely. This motional mechanism has indeed been shown to lead to ^2H NMR lines as narrow as 0.7 kHz in a low molar mass crystalline solid,¹⁵ provided the jumps occur between all four sites.

The extent to which the mobility at the two labeled sites differs is evident upon inspection of Figure 3. At temperatures below 300 K or above 410 K, the values of Φ_{ma} at the two sites are indistinguishable. At temperatures in between these two extremes, however, Φ_{ma} is consistently larger for the outer-labeled sample than for the sample labeled at the center of the hard segment, indicating greater mobility at the surface of the crystal than at its center and that the onset of mobility within the crystals begins at the surface and proceeds inward. While the absolute value of Φ_{ma} will be affected by the degree of phase separation, the absolute difference of Φ_{ma} for the two samples is unaffected, so that it is related to the difference in the extent to which the motions are constrained, and has a maximum value of 0.27 ± 0.04 at 380 K, well below $T_m(\text{hard})$.

Dynamic coupling between the chains in the hard and soft domains is known to occur in block copolymers^{17,18} so that the more rigid hard domains would restrict the mobility of the soft segments,^{10,19} which in turn induce motions within the hard domains, especially at the interface and at elevated temperatures. Comparison of Φ_{ma} of the multi- and triblock polyurethanes suggests that the dynamic coupling between the two phases occurs to a significant extent. The value of Φ_{ma} for the central-labeled multiblock polymer is consistently below that of the equivalent triblock polymer at $T > 350$ K, and the plateau region in Φ_{ma} observed in the case of the multiblock copolymers is obscured in the case of the triblock polyurethanes (see Figure 3a). This suggests that there is greater dynamic coupling in the triblock copolymer. Such a phenomenon can be attributed to the difference in mobility of the soft segments comprising the two classes of polymers. Both ends of the soft segments are covalently bound to the hard segments in the multiblocks, whereas in the triblocks, only one end is linked to the hard segment while the other is free. The additional freedom permits the soft segment to induce motions within the hard domain of the triblock more easily. The effect is only observed once the ^2H NMR spectra of the hard-segment oligomers indicate the onset of motions at the center of the crystal, i.e., ca. 360 K, so that the onset of mobility within the ordered crystalline matrix facilitates the dynamic coupling with the more mobile soft segments.

Below 350 K, the central-labeled oligomer does not exhibit a motionally narrowed peak (see Figure 3b), so the presence of such a peak in the polymers can be attributed to incomplete phase separation. By comparison of Φ_{ma} of a non-phase-separating model system, i.e., a polymer where the soft segments are chain extended by individual labeled piperazine rings, and hard-segment oligomers, it was shown that the value of Φ_{ma} at temperatures ranging from 310 to 350 K for the polymer labeled in the central ring is equal to the fraction of hard segments dispersed in the soft phase.⁹ For the triblock polyurethanes, this would correspond to a value of Φ_{ma} of 0.09 ± 0.02 . This upper limit seems reasonable in light of the value reported by Kornfield et al., 0.15 ± 0.02 , for the corresponding multiblock polyurethanes.

Conclusions

The motional behavior within the hard domain of the polyurethane triblock polymers can be divided into different temperature regimes. At $T < 300$ K, the soft phase has solidified to some extent and hinders the mobility of that part of the hard segments at the interface. Between 310 and 350 K the dual character of the ^2H NMR spectra is indicative of the heterogeneous nature of the system, and Φ_{ma} of the center-labeled sample indicates

that $9 \pm 2\%$ of the hard segments are "dissolved" in the soft phase. Dynamic coupling between the two types of segments comprising the polymer is observed and occurs to an even greater extent between 360 and 400 K. Above 360 K, tetrahedral jumps can occur at a sufficient rate within the ordered hard phase so that the spectra are motionally averaged. Last, at $T > 440$ K, the hard segments are disordered.

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References and Notes

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